

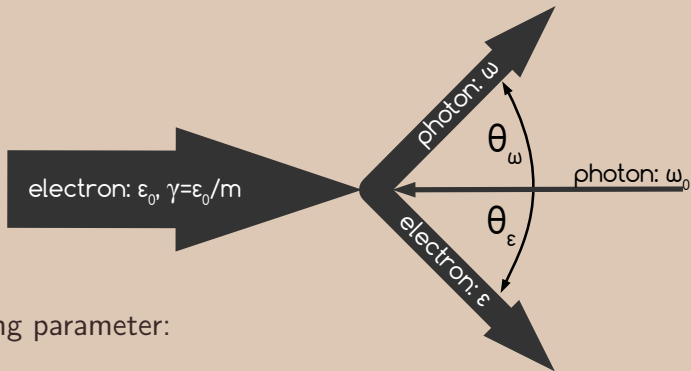
Compton polarimetry by scattered electrons (translated to english)

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Inverse Compton scattering: kinematics

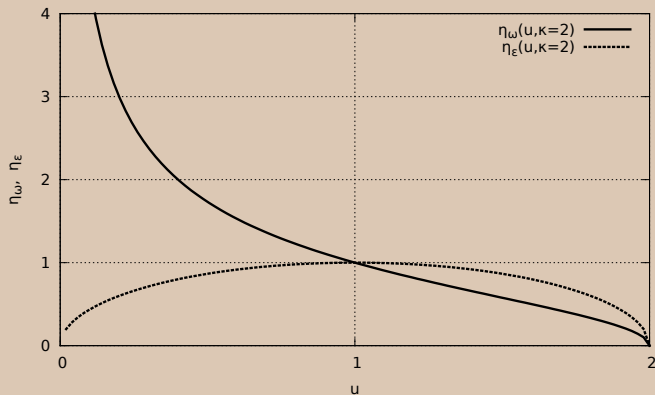


Scattering parameter:

$$u = \frac{\omega}{\epsilon} = \frac{\theta_\epsilon}{\theta_\omega} = \frac{\omega}{\epsilon_0 - \omega} = \frac{\epsilon_0 - \epsilon}{\epsilon}; \quad u \in [0, \kappa], \quad \text{where } \kappa = \frac{4\omega_0\epsilon_0}{m^2}.$$

$$\text{Scattering angles: } \eta_\omega \equiv \gamma\theta_\omega = \sqrt{\frac{\kappa}{u} - 1}; \quad \eta_\epsilon \equiv \gamma\theta_\epsilon = u\sqrt{\frac{\kappa}{u} - 1}.$$

Inverse Compton scattering: kinematics



$\kappa = 2$ when, e. g., $\omega_0 = 2.33$ eV and $\varepsilon_0 = 56$ GeV.

For any κ , if $u = \frac{\kappa}{2}$ one has $\eta_\omega = 1$, $\eta_\epsilon = \frac{\kappa}{2}$, i. e. $\max(\theta_\epsilon) = \frac{2\omega_0}{m}$.

Maximum electron scattering angle does not depend on beam energy!

Inverse Compton scattering: cross section

$$d\sigma = \left\{ \frac{1}{\kappa(1+u)^2} \left(2 + \frac{u^2}{1+u} + 4 \frac{u}{\kappa} \left[\frac{u}{\kappa} - 1 \right] \left[1 - \xi_{\perp} \cos(2(\varphi - \varphi_{\perp})) \right] \right) + \xi_{\circ} \left(\zeta_{\parallel} \frac{u(u+2)(\kappa-2u)}{\kappa^2(1+u)^3} - \zeta_{\perp} \frac{2u^2 \sqrt{\kappa/u-1}}{\kappa^2(1+u)^3} \sin \varphi \right) \right\} r_e^2 d\varphi du,$$

Modified Stokes parameters:

ξ_{\perp} и φ_{\perp} – degree and direction of laser linear polarization,

ξ_{\circ} – degree of laser circular polarization,

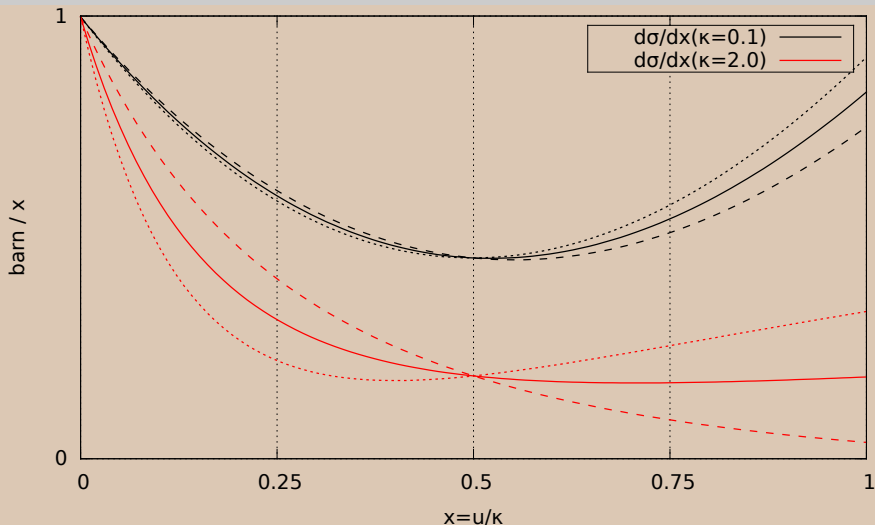
ζ_{\parallel} – longitudinal electron beam polarization degree,

ζ_{\perp} – transverse vertical electron beam polarization degree.

Laser beam: $\sqrt{\xi_{\perp}^2 + \xi_{\circ}^2} = 1$, $\xi_{\perp} \in [0, 1]$, $\xi_{\circ} \in [-1, 1]$.

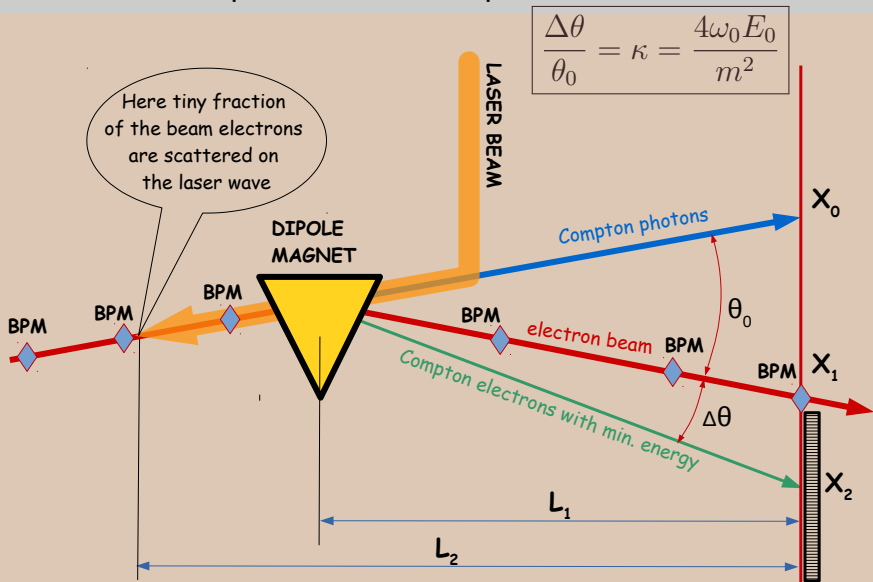
Electron beam: $\sqrt{\zeta_{\perp}^2 + \zeta_{\parallel}^2} < 1$, $\zeta_{\perp} \in [-1, 1]$, $\zeta_{\parallel} \in [-1, 1]$.

Inverse Compton scattering: cross section



Dashed lines illustrate the influence of $\xi_0 \zeta_{||}$

General concept of I.C.S. experiments



Compton polarimeters for electron beam

VEPP-2, VEPP-4, LEP, HERA, SLC ... ILC, FCC ...

At first Compton polarimeters it was general to deal with backscattered photons.

At higher electron energies the divergence of γ -beam is small, high energy SR photons appear, etc., it seems reasonable to look at scattered electrons.

Maximum scattering angle of electron does not depend on its initial energy: $\max(\theta_\epsilon) = 2\omega_0/m$.

If $\omega_0 = 2.33$ eV one has $\max(\theta_\epsilon) \simeq 10 \mu\text{rad}$.

Vertical beam energy spread: $\sigma'_y = \sqrt{\epsilon_y/\beta_y} \ll \max(\theta_\epsilon)$.

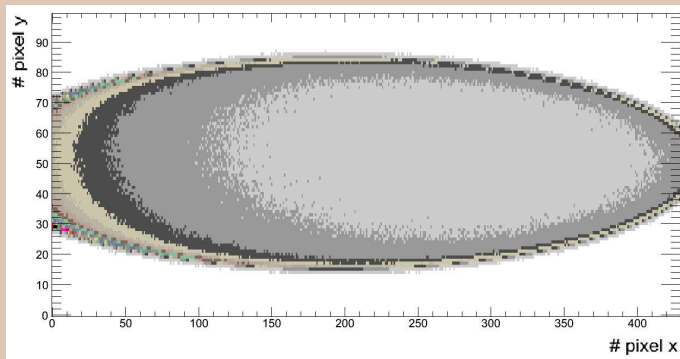
An example: $\epsilon_y = 100$ pm и $\beta_y = 100$ m gives $\sigma'_y = 1 \mu\text{rad}$.

ILC note: LC-M-2012-001

A Transverse Polarimeter for a Linear Collider of 250 GeV e Beam Energy

Itai Ben Mordechai and Gideon Alexander

“... For the detection of the scattered electrons we consider only a position measurement using a Silicon pixel detector placed at a distance of 37.95 m from the Compton IP. The active dimension of the detector is $2 \times 200 \text{ mm}^2$. The size of the pixels cell taken is $50 \times 400 \mu\text{m}^2$ similar to the one used in the ATLAS detector [9]. This scheme yields an approximate two dimensional resolution of $14.4 \times 115.5 \mu\text{m}^2$ [10] with a data read-out rate of ...”



Angles of scattered electron after bending dipole

The energy of scattered electron depends on u as: $\varepsilon = \varepsilon_0/(1 + u)$.
Such an electron will be bent in a dipole to an angle:

$$\theta_s = \frac{A}{\varepsilon} = \frac{A}{\gamma m}(1 + u); \quad \eta_s \equiv \gamma \theta_s = \eta_0 + u\eta_0,$$

where A is the dipole “force”, $\gamma = \varepsilon_0/m$, $\eta_0 = A/m$.

Introducing new parameters η_x, η_y :

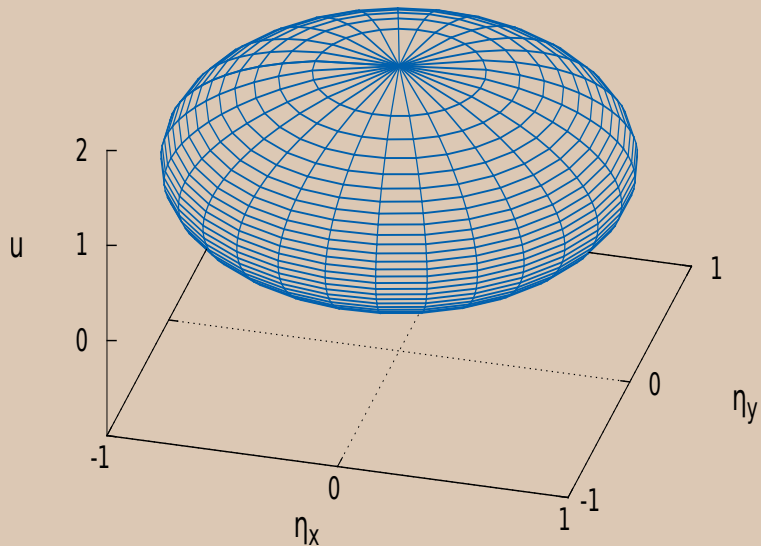
$$\begin{cases} \eta_x \equiv \eta_s - \eta_0 = u\eta_0 + u\sqrt{\kappa/u - 1} \cos \varphi \\ \eta_y = u\sqrt{\kappa/u - 1} \sin \varphi \end{cases}$$

one gets the square equation for u :

$$(\eta_x - u\eta_0)^2 + \eta_y^2 = u(\kappa - u).$$

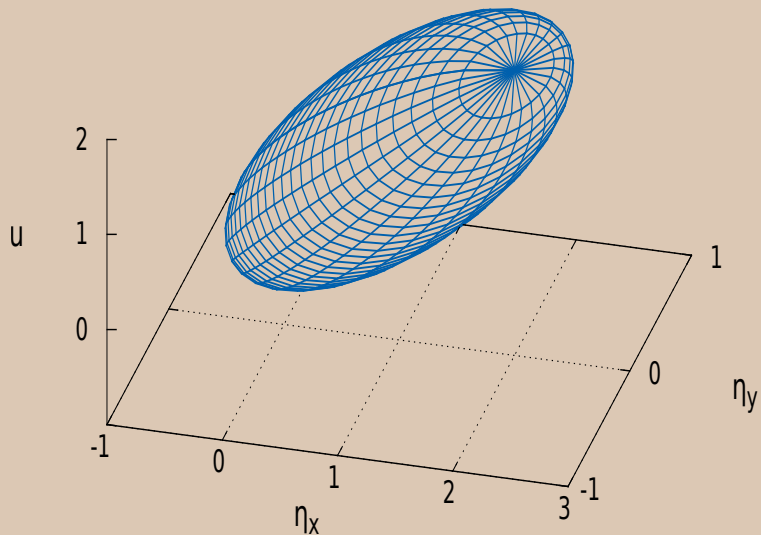
Scattering surface

$$\kappa=2, \eta_0=0$$



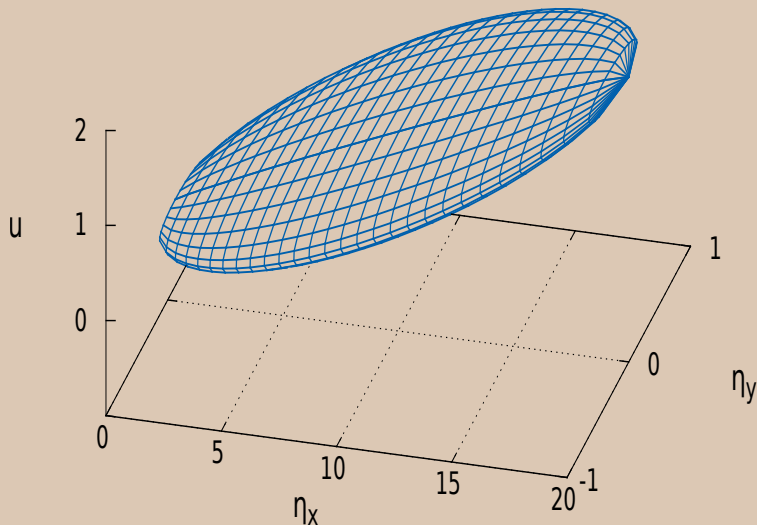
Scattering surface

$$\kappa=2, \eta_0=1$$



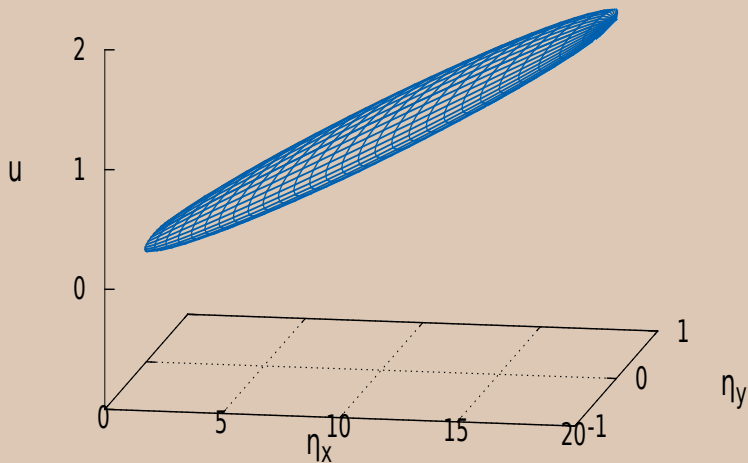
Scattering surface

$$\kappa=2, \eta_0=10$$



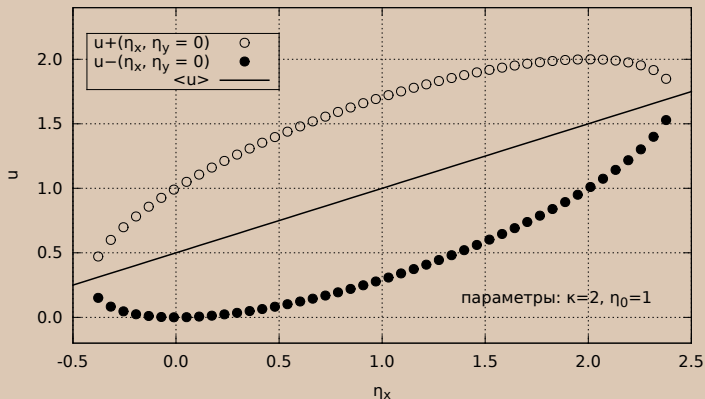
Scattering surface

$$\kappa=2, \eta_0=10$$



Solving the equation:

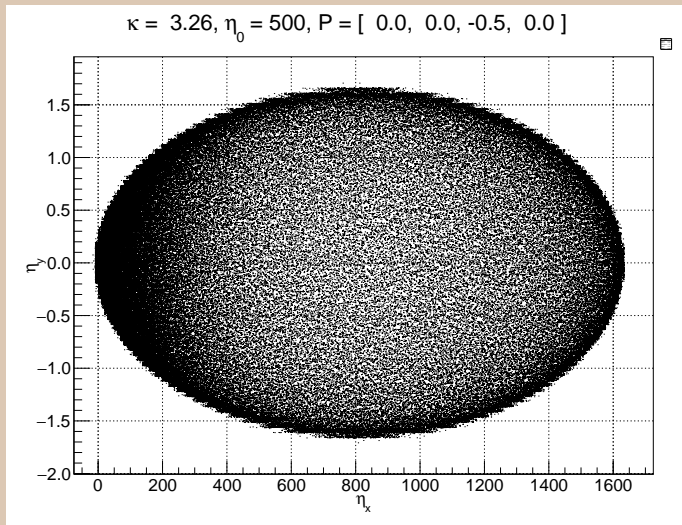
$$u = \frac{\kappa/2 + \eta_0\eta_x \pm 2J}{1 + \eta_0^2}, \text{ where } J = \sqrt{\kappa^2/4 - \eta_x^2 - \eta_y^2(1 + \eta_0^2) + \kappa\eta_0\eta_x}$$



$$\langle u \rangle = \frac{\kappa/2 + \eta_0\eta_x}{1 + \eta_0^2} \xrightarrow{\eta_0 \gg 1} \frac{\eta_x}{\eta_0}$$

And also: $d\varphi du = d\eta_x d\eta_y / J$

Ellipse in η_x, η_y plane



Sizes of ellipse are [rad]: $O_y = \frac{4\omega_0}{m}, O_x = \frac{4\omega_0}{m} \sqrt{1 + (\gamma\theta_0)^2}$

Cross section: $(u, \varphi) \rightarrow (\eta_x, \eta_y)$

$$\begin{aligned}
 d\sigma &= \left\{ \frac{1}{\kappa(1+u)^2} \left(2 + \frac{u^2}{1+u} + 4 \frac{u}{\kappa} \left[\frac{u}{\kappa} - 1 \right] \left[1 - \xi_{\perp} \cos(2(\varphi - \varphi_{\perp})) \right] \right) + \right. \\
 &\quad \left. + \xi_{\odot} \left(\zeta_{\parallel} \frac{u(u+2)(\kappa-2u)}{\kappa^2(1+u)^3} - \zeta_{\perp} \frac{2u^2 \sqrt{\kappa/u - 1}}{\kappa^2(1+u)^3} \sin \varphi \right) \right\} r_e^2 d\varphi du = \\
 &= \left\{ \frac{1}{\kappa(1+u)^2} \left(2 + \frac{u^2}{1+u} + 4 \frac{u}{\kappa} \left[\frac{u}{\kappa} - 1 \right] \left[1 - \xi_{\perp} \left(1 - \frac{2\eta_y^2 \cos^2 \varphi_{\perp}}{u(\kappa - u)} \right) \right] \right) + \right. \\
 &\quad \left. + \xi_{\odot} \left(\zeta_{\parallel} \frac{u(u+2)(\kappa-2u)}{\kappa^2(1+u)^3} - \zeta_{\perp} \frac{2u\eta_y}{\kappa^2(1+u)^3} \right) \right\} \frac{r_e^2 d\eta_x d\eta_y}{J(\eta_x, \eta_y)} \times 2.
 \end{aligned}$$

Cross section: $(\eta_x, \eta_y) \rightarrow (r, \phi)$

$$d\sigma = \frac{\Sigma(\eta_x, \eta_y) d\eta_x d\eta_y}{\sqrt{\kappa^2/4 - \eta_x^2 - \eta_y^2(1 + \eta_0^2) + \kappa\eta_0\eta_x}}$$

Ellipse \rightarrow circle: $(x = \frac{\eta_x - \kappa\eta_0/2}{\sqrt{1 + \eta_0^2}}; y = \eta_y)$, then:

$$J(x, y) = \sqrt{(1 + \eta_0^2)(\kappa^2/4 - x^2 - y^2)} = \sqrt{(1 + \eta_0^2)(\kappa^2/4 - r^2)},$$

$$d\eta_x d\eta_y = \sqrt{1 + \eta_0^2} dx dy = \sqrt{1 + \eta_0^2} r dr d\phi,$$

$$\langle u \rangle = \frac{\kappa}{2} + \frac{x\eta_0}{\sqrt{1 + \eta_0^2}} = \frac{\kappa}{2} + \frac{r \cos \phi \cdot \eta_0}{\sqrt{1 + \eta_0^2}}.$$

Finally:
$$d\sigma = \frac{\Sigma(r, \phi) r dr d\phi}{\sqrt{\kappa^2/4 - r^2}}$$

Take emittance into account: $f(\eta_x, \eta_y) =$

$$= \int_0^R \frac{r dr}{\sqrt{R^2 - r^2}} \int_0^{2\pi} \Sigma(r, \phi) \exp \left[-\frac{(r \cos \phi - \eta_x^*)^2}{2\sigma_x^{*2}} - \frac{(r \sin \phi - \eta_y)^2}{2\sigma_y^2} \right] d\phi,$$

where $R = \kappa/2$, $\eta_x^* = (\eta_x - \kappa\eta_0/2)/\sqrt{1 + \eta_0^2}$, $\sigma_x^* = \sigma_x/\sqrt{1 + \eta_0^2}$.

We divide the interval $[0 : R]$ into N identical segments:

$$r_{i,B} = R \cdot i/N \equiv RB_i, \quad r_{i,E} = R(i+1)/N \equiv RE_i.$$

$$\int_{r_{i,B}}^{r_{i,E}} \frac{r dr}{\sqrt{R^2 - r^2}} = R \mathcal{I}_i; \quad \int_{r_{i,B}}^{r_{i,E}} \frac{r^2 dr}{\sqrt{R^2 - r^2}} = R^2 \mathcal{J}_i$$

Weighted average radius of i ring: $\bar{r}_i = R \cdot \mathcal{J}_i/\mathcal{I}_i$.

The sum instead of integral: $f(\eta_x, \eta_y) =$

$$= R \sum_{i=0}^{N-1} \mathcal{I}_i \Delta\phi_i \sum_{j=0}^{\lfloor 2\pi(i+1) \rfloor} \Sigma_{ij} \exp \left[-\frac{(Rx_{ij} - \eta_x^*)^2}{2\sigma_x^{*2}} - \frac{(Ry_{ij} - \eta_y)^2}{2\sigma_y^2} \right],$$

where

$$\Delta\phi_i = 2\pi / \lfloor 2\pi(i+1) \rfloor$$

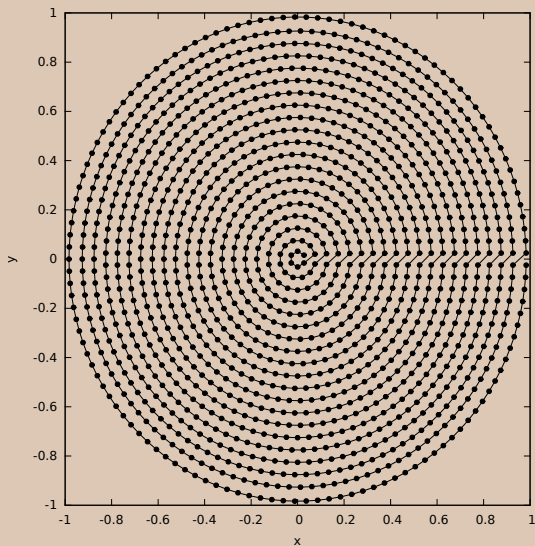
$$x_{ij} = \frac{\mathcal{J}_i}{\mathcal{I}_i} \cdot \frac{\sin((j+1)\Delta\phi_i) - \sin(j\Delta\phi_i)}{\Delta\phi_i}$$

$$y_{ij} = \frac{\mathcal{J}_i}{\mathcal{I}_i} \cdot \frac{\cos(j\Delta\phi_i) - \cos((j+1)\Delta\phi_i)}{\Delta\phi_i}$$

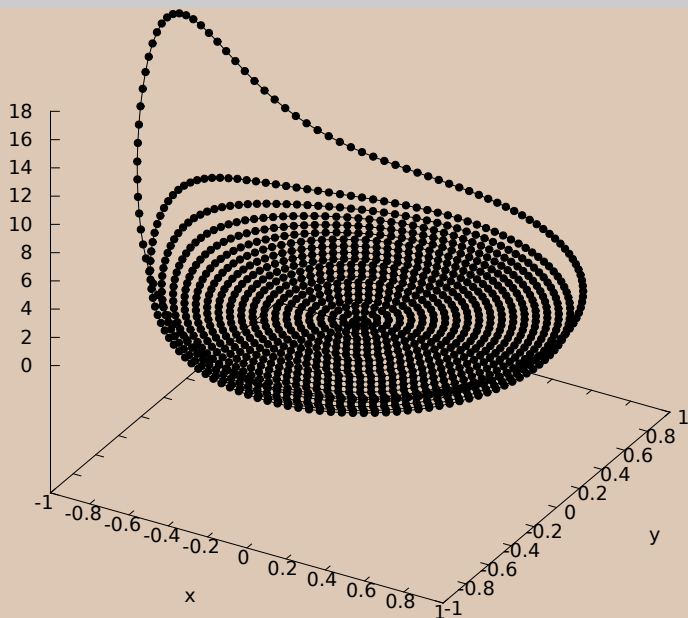
$$u_{ij} = R(1 + x_{ij}) \in [0 : 2R]$$

$$\Sigma_{ij} = \frac{(1 + x_{ij}^2)(1 + u_{ij}) + u_{ij}^2}{\kappa(1 + u_{ij})^3} - P_{\parallel} \frac{x_{ij}u_{ij}(2 + u_{ij})}{\kappa(1 + u_{ij})^3} - P_{\perp} \frac{y_{ij}(1 + x_{ij})}{2(1 + u_{ij})^3}$$

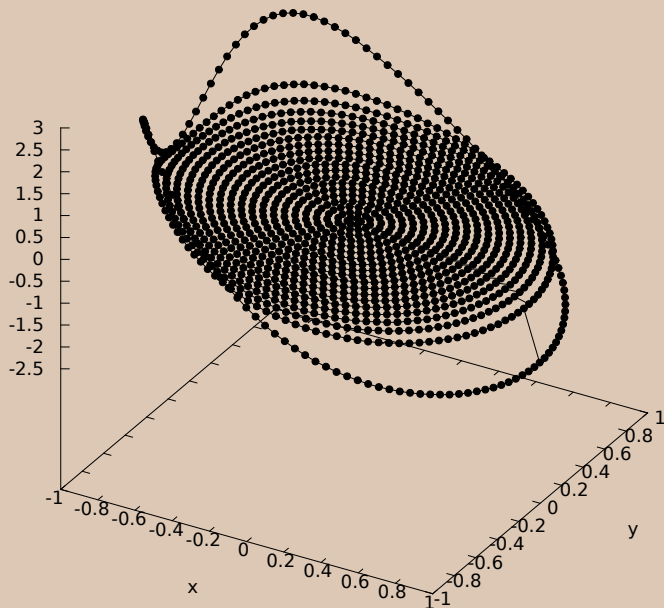
x_{ij} VS y_{ij}



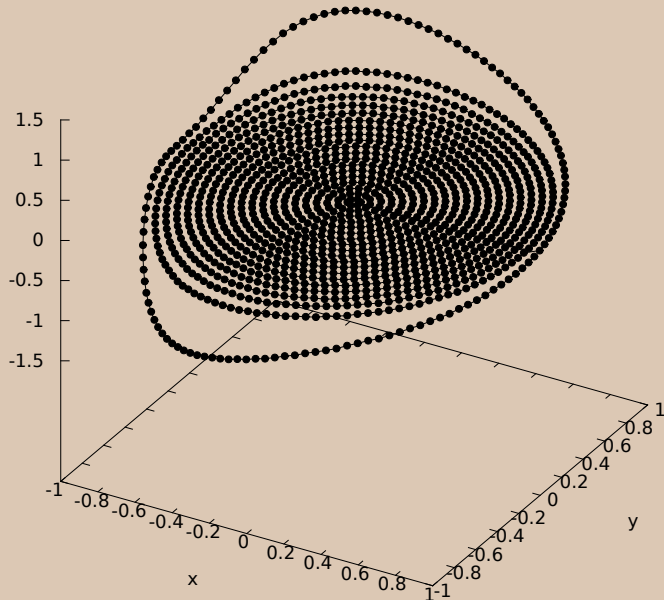
Unpolarized cross section



Contribution of 100% longitudinal polarization



Contribution of 100% transverse polarization



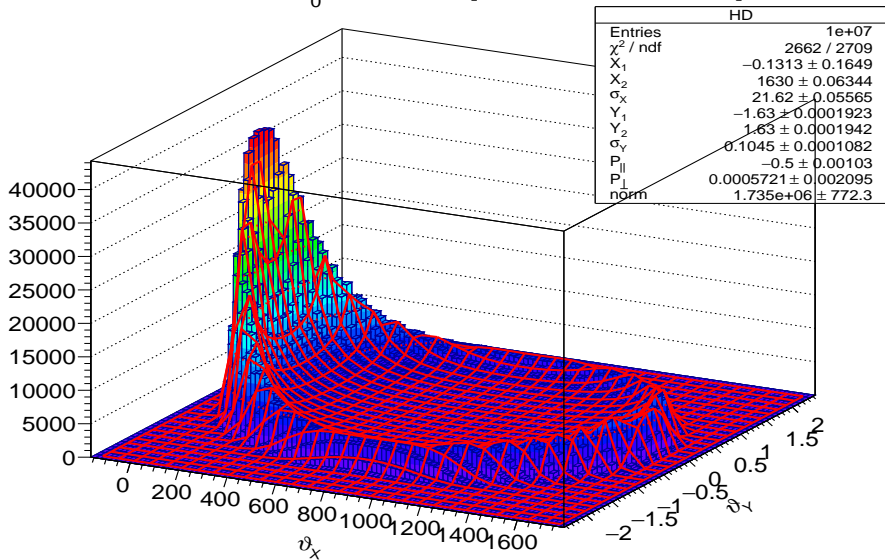
The function to fit the η_x, η_y distribution of scattered electrons:

$$f(X, Y) = I \cdot \sum_{i,j} \frac{\mathcal{I}_i \Delta\varphi_i \Sigma_{ij}(P_{\parallel}, P_{\perp})}{2\pi\sigma_X\sigma_Y} \times \\ \times \exp \left[-\frac{\left(\frac{X_2 - X_1}{2} \bar{x}_{ij} + \frac{X_1 + X_2}{2} - X \right)^2}{2\sigma_X^2} \right] \times \\ \times \exp \left[-\frac{\left(\frac{Y_2 - Y_1}{2} \bar{y}_{ij} + \frac{Y_1 + Y_2}{2} - Y \right)^2}{2\sigma_Y^2} \right].$$

Fit parameters: $X_1, X_2, \sigma_X, Y_1, Y_2, \sigma_Y, P_{\parallel}, P_{\perp}, I$ (9 at all)

Fit results: longitudinal polarization 50%

$$\kappa = 3.26, \vartheta_0 = 500, P = [0.0, 0.0, -0.5, 0.0]$$



Fit results: longitudinal polarization 50%

Fitting range along X – from 200

PARAMETER CORRELATION COEFFICIENTS

NO.	GLOBAL	1	2	3	4	5	6	7	8	9
1	0.63806	1.000	-0.311	0.114	-0.332	0.339	-0.009	-0.360	0.018	-0.465
2	0.61959	-0.311	1.000	-0.541	0.269	-0.282	0.123	0.219	-0.013	0.145
3	0.57955	0.114	-0.541	1.000	-0.193	0.188	-0.262	-0.112	0.004	-0.052
4	0.57920	-0.332	0.269	-0.193	1.000	-0.468	0.418	0.146	0.010	0.144
5	0.60208	0.339	-0.282	0.188	-0.468	1.000	-0.443	-0.146	0.024	-0.166
6	0.57481	-0.009	0.123	-0.262	0.418	-0.443	1.000	0.013	-0.005	0.022
7	0.38479	-0.360	0.219	-0.112	0.146	-0.146	0.013	1.000	-0.003	0.231
8	0.03962	0.018	-0.013	0.004	0.010	0.024	-0.005	-0.003	1.000	-0.003
9	0.47128	-0.465	0.145	-0.052	0.144	-0.166	0.022	0.231	-0.003	1.000

FCN=2662.5 FROM MIGRAD STATUS=CONVERGED 257 CALLS 258 TOTAL

EDM=3.93455e-08 STRATEGY= 1 ERROR MATRIX UNCERTAINTY 0.8 per cent

EXT PARAMETER

PARABOLIC

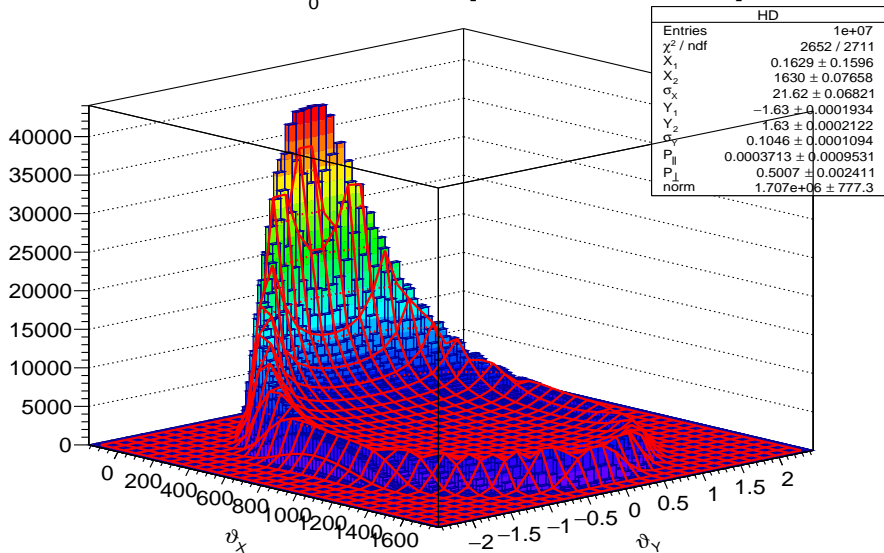
MINOS ERRORS

NO.	NAME	VALUE	ERROR	NEGATIVE	POSITIVE
1	X_{1}	-1.31308e-01	1.64882e-01		
2	X_{2}	1.62998e+03	6.34381e-02		
3	#sigma_{X}	2.16201e+01	5.56481e-02		
4	Y_{1}	-1.62981e+00	1.92272e-04		
5	Y_{2}	1.62973e+00	1.94174e-04		
6	#sigma_{Y}	1.04485e-01	1.08179e-04		
7	P_{#paral}	-5.00034e-01	1.02951e-03		
8	P_{#perp}	5.72060e-04	2.09542e-03		
9	norm	1.73486e+06	7.72345e+02		

2DFit : Real Time = 140.38 seconds Cpu Time = 156.49 seconds

Fit results: transverse polarization 50%

$$\kappa = 3.26, \vartheta_0 = 500, P = [0.0, 0.0, 0.0, 0.5]$$



Fit results: transverse polarization 50%

Fitting range along X – from 200

PARAMETER CORRELATION COEFFICIENTS

NO.	GLOBAL	1	2	3	4	5	6	7	8	9
1	0.69373	1.000	-0.318	0.138	-0.415	0.375	-0.015	-0.320	0.013	-0.537
2	0.62430	-0.318	1.000	-0.556	0.270	-0.250	0.125	0.232	-0.003	0.154
3	0.59633	0.138	-0.556	1.000	-0.185	0.189	-0.274	-0.145	0.012	-0.071
4	0.62861	-0.415	0.270	-0.185	1.000	-0.432	0.432	0.174	0.095	0.206
5	0.58723	0.375	-0.250	0.189	-0.432	1.000	-0.397	-0.137	0.153	-0.210
6	0.58308	-0.015	0.125	-0.274	0.432	-0.397	1.000	0.022	-0.008	0.017
7	0.36220	-0.320	0.232	-0.145	0.174	-0.137	0.022	1.000	0.011	0.243
8	0.23947	0.013	-0.003	0.012	0.095	0.153	-0.008	0.011	1.000	0.022
9	0.54496	-0.537	0.154	-0.071	0.206	-0.210	0.017	0.243	0.022	1.000

FCN=2651.75 FROM MIGRAD STATUS=CONVERGED 258 CALLS 259 TOTAL
EDM=4.09631e-07 STRATEGY= 1 ERROR MATRIX UNCERTAINTY 0.4 per cent

EXT	PARAMETER	PARABOLIC	MINOS ERRORS		
NO.	NAME	VALUE	ERROR	NEGATIVE	POSITIVE
1	X_{1}	1.62941e-01	1.59586e-01		
2	X_{2}	1.63002e+03	7.65815e-02		
3	#sigma_{X}	2.16220e+01	6.82096e-02		
4	Y_{1}	-1.62982e+00	1.93423e-04		
5	Y_{2}	1.63003e+00	2.12161e-04		
6	#sigma_{Y}	1.04595e-01	1.09394e-04		
7	P_{#paral}	3.71312e-04	9.53123e-04		
8	P_{#perp}	5.00724e-01	2.41133e-03		
9	norm	1.70728e+06	7.77293e+02		

2DFit : Real Time = 131.11 seconds Cpu Time = 158.11 seconds

Results & discussion

- Analysis of scattered electrons (SE) properties was performed.
- $X - Y$ distribution of SE after a bending magnet looks like an ellipse with sharp edges, defined by scattering kinematics.
- Vertical size of this ellipse is independent on the beam energy. Low electron beam emittance (vertical) is required to observe the original cross section.
- Horizontal size of the ellipse depends on the product of two parameters, $\gamma\theta_0$, where $\gamma = \varepsilon_0/m$ and θ_0 is the beam bending angle in the same dipole. Thus $\gamma\theta_0$ can be measured with a relative statistical accuracy of at least 10^{-4} .
- Either transverse or longitudinal beam polarization can be evaluated from the same measurement.

Thank You!